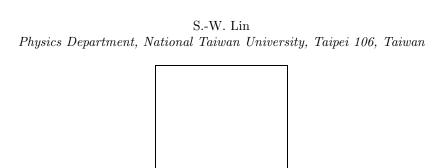
$B \rightarrow$ Baryon decays in Belle



We report recent observations of baryonic B decays with charmless and charmed baryons in the final state. We show the angular distributions of the di-baryon low-mass enhancements in the charmless three-body baryonic B decays and the branching fractions of B decays with two charmed baryons or charmonium in the final states. We also report the observation of the decay $\eta_c \to \Lambda \bar{\Lambda}$ at Belle.

1 Introduction

Observations of several baryonic B decays have been reported by Belle. The measured branching fractions for charmless and charmed baryonic B decays are shown in Fig. 1. In the charmless final states, only the three-body decays have been observed 1,2,3,4,5. In this contribution, we report on the angular distribution of the di-baryon low-mass enhancements seen in the charmless three-body baryonic B decays. The data support the quark fragmentation interpretation, while the gluonic resonance picture is disfavored. In the charmed final states, we observed the two-body and three-body decays with two charmed baryons or charmonium. From the latter we can extract the branching fractions of η_c into baryon pairs $p\bar{p}$ and - for the first time - into $\Lambda\bar{\Lambda}$. Measuring decay rates of η_c to different di-baryon modes is a very useful check for theoretical predictions and can shed light on quark-diquark dynamics. The data sample was collected with Belle detector at the KEKB asymmetric-energy $e^+e^-(3.5 \text{ GeV})$ on 8.0 GeV) collider KEKB operates at the $\Upsilon(4S)$ resonance ($\sqrt{s}=10.58 \text{ GeV}$) with a peak luminosity that has exceeded $1.5 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$.

2 Charmless baryonic B decays

Observations of several charmless baryonic B decays have been reported at Belle. One common feature of these observations is the peaking of the di-baryon mass spectra toward threshold.

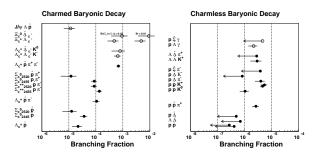


Figure 1: Branching fractions for charmless and charmed baryonic B decays. The open circles indicate new measurements from Belle.

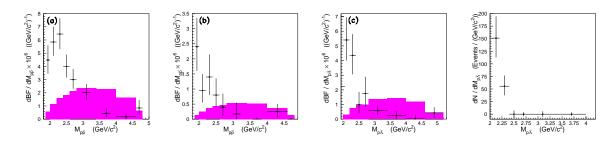


Figure 2: Differential branching fraction for (a) $p\bar{p}K^+$, (b) $p\bar{p}K^0_S$, (c) $p\bar{\Lambda}\pi^-$ and (d) $p\bar{\Lambda}\gamma$ modes as a function of di-baryon pair mass. The shaded distribution shows the expectation from a phase-space MC simulation with area scaled to the signal yield. A charmonium veto has been applied in (a) and (b).

We first measured the differential branching fractions for $(a)B^+ \to p\bar{p}K^+$, $(b)B^0 \to p\bar{p}K^0_S$, $(c)B^0 \to p\bar{\Lambda}\pi^-$ and $(d)B^+ \to p\bar{\Lambda}\gamma$ modes. There are two kinematic variables in the center of mass frame which are usually used to extract the B candidates: the beam energy constrained mass $M_{\rm bc} = \sqrt{E_{\rm beam} - p_B^2}$, and the energy difference $\Delta E = E_B - E_{\rm beam}$, where $E_{\rm beam}$ is the beam energy, and p_B and E_B are the momentum and energy of B candidates. B signal yield is obtained by a 2D fit to the $(M_{\rm bc}, \Delta E)$ distribution for each bin of di-baryon invariant mass. The efficiency as a function of the di-baryon mass is based on the MC simulation. The differential branching fraction for each bin in the di-baryon mass is obtained from the fitted B yield and the signal efficiency (Fig. 2). The branching fraction is a sum of the differential branching fractions^{4,5}.

The angular distribution of the proton is studied in the di-baryon system. Fig. 3 shows the angular distributions for the four decay channels. The angle θ_p is defined for (a) as the angle between the directions of \bar{p} and K^+ in the $p\bar{p}$ rest frame. There is a clear forward peak and the angular asymmetry, defined as

$$A = \frac{N_{+} - N_{-}}{N_{+} + N_{-}} \tag{1}$$

where N_+ and N_- stand for the efficiency corrected B yield with $\cos \theta_p > 0$ and $\cos \theta_p < 0$, respectively, amounts to $0.59^{+0.08}_{-0.07}$ for the $p\bar{p}K^+$ mode. The asymmetry of the distribution indicates that the fragmentation picture is favored. Antiprotons are emitted along the K^+ direction most of the time, which can be explained by a parent $\bar{b} \to \bar{s}$ penguin transition followed by $\bar{s}u$ fragmentation into the final state. The $\cos \theta_p$ distribution for (b) is flat, but we have to note that in this case the statistics is low and K_S^0 carries no flavor information. If in (c) we choose the θ_p angle of p with respect to the π^- in the rest frame of the $\bar{\Lambda}p$ system, the $\cos \theta_p$ distribution is quite flat. However, the $\cos \theta_p$ shows a forward peak structure if the angle is defined as the p direction relative to the $\bar{\Lambda}$ direction in the $p\pi^-$ rest frame. It is evident that the fragmentation interpretation is supported: the proton tends to emerge parallel to the

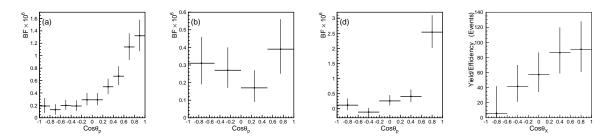


Figure 3: Branching fraction vs. $\cos \theta_p$ in the di-baryon system for (a) $p\bar{p}K^+$, (b) $p\bar{p}K^0_S$, (c) $p\bar{\Lambda}\pi^-$ and (d) $p\bar{\Lambda}\gamma$.

 $\bar{\Lambda}$ baryon. The angle θ_X of (d) is measured between the proton direction and the γ direction in the baryon pair rest frame. There is also a clear forward structure in the $\cos \theta_X$ distribution and the angular asymmetry A is $0.36^{+0.23}_{-0.20}$. This distribution supports the $b \to s \gamma$ fragmentation picture where the Λ tends to emerge opposite to the direction of the photon⁵.

3 Charmed baryonic B decays

In the $B^{+/0}\to \bar\Xi_c^{0/-}\Lambda_c^+$ analysis⁸, we reconstruct the following decay modes: $\Xi_c^0\to\Xi^-\pi^+$ and $\Lambda K^-\pi^+$, $\Xi_c^+\to\Xi^-\pi^+\pi^+$, $\Lambda_c^+\to pK^-\pi^+$, $\Xi^-\to\Lambda\pi^-$ and $\Lambda\to p\pi^-$. We use a simultaneous two-dimensional binned maximum likelihood fit to the ΔE vs. $M_{\rm bc}$ distributions (for the two Ξ_c^0 channels) with a common value of $\mathcal{B}(B^+\to\bar\Xi_c^0\Lambda_c^+)\times\mathcal{B}(\bar\Xi_c^0\to\bar\Xi^+\pi^-)\times\mathcal{B}(\Lambda_c^+\to pK^-\pi^+)$. For this fit, we constrain the ratio $\mathcal{B}(\Xi_c^0\to\Lambda K^-\pi^+)/\mathcal{B}(\Xi_c^0\to\Xi^-\pi^+)$ to the recent Belle measurement of $1.07\pm0.12\pm0.07$. The fit gives the product of branching fractions of $\mathcal{B}(B^+\to\bar\Xi_c^0\Lambda_c^+)\times\mathcal{B}(\bar\Xi_c^0\to\bar\Xi^+\pi^-)$ with the value of $4.8_{-0.9}^{+1.0}\pm1.1\pm1.2$ and a statistical significance of $8.7\ \sigma$. We also check the $B^0\to\bar\Xi_c^-\Lambda_c^+$ mode which is an isospin partner of the $B^+\to\bar\Xi_c^0\Lambda_c^+$ mode. The product of branching fraction of $\mathcal{B}(B^0\to\bar\Xi_c^-\Lambda_c^+)\times\mathcal{B}(\bar\Xi_c^-\to\bar\Xi^+\pi^-\pi^-)$ is measured to be $(9.3_{-2.8}^{+3.7}\pm1.9\pm2.4)\times10^{-5}$. The uncertainties in the products of branching ratios are statistical, systematic and the uncertainty from the $\Lambda_c^+\to pK^-\pi^+$ branching fraction.

The $B^+ \to \Lambda_c^+ \Lambda_c^- K^+$ and $B^0 \to \Lambda_c^+ \Lambda_c^- K^0$ decays are three-body decays that proceed via a $b \to c\bar{c}s$ transition. We detect the Λ_c^+ via the $\Lambda_c^+ \to pK^-\pi^+$, $p\bar{K}^0$ and $\Lambda\pi^+$ decay channels. When a Λ_c^+ and Λ_c^- are combined as B decay daughters, at least one of Λ_c^\pm is required to have been reconstructed via the $pK^\mp\pi^\pm$ decay process. Here, the parameter mass difference ΔM_B is used instead of the energy difference ΔE , since ΔE shows a correlation with $M_{\rm bc}$. The mass difference is defined as $\Delta M_B \equiv M(B) - m_B$, where M(B) is the reconstructed mass of the B candidate and m_B is the world average B meson mass. Fig. 4 shows ΔM_B and $M_{\rm bc}$ projections for $B^+ \to \Lambda_c^+ \Lambda_c^- K^+$ and $B^0 \to \Lambda_c^+ \Lambda_c^- K^0$ decays. A two-dimensional binned maximum likelihood fit is performed to determine the signal yield. From the fit we obtain signal yields of $48.5_{-6.8}^{+7.5}$ and $10.5_{-3.1}^{+3.8}$ events with statistical significances of 15.4 σ and 6.6 σ , for $B^+ \to \Lambda_c^+ \Lambda_c^- K^+$ and $B^0 \to \Lambda_c^+ \Lambda_c^- K^0$, respectively. Hence, we obtain the branching fractions of

$$\mathcal{B}(B^+ \to \Lambda_c^+ \Lambda_c^- K^+) = (6.5^{+1.0}_{-0.9} \pm 1.1 \pm 3.4) \times 10^{-4}$$

$$\mathcal{B}(B^0 \to \Lambda_c^+ \Lambda_c^- K^0) = (7.9^{+2.9}_{-2.3} \pm 1.2 \pm 4.1) \times 10^{-4}$$

where the first and the second errors are statistical and systematic, respectively. The last error is due to the 52 % uncertainty in the absolute branching fraction, $\mathcal{B}(\Lambda_c^+ \to pK^-\pi^+)$.

We study two-body baryonic decays of charmonia in the B meson decays, $B^+ \to p\bar{p}K^+$ and $B^+ \to \Lambda\bar{\Lambda}K^+$. The B signal yields are obtained from 10 MeV/c² wide $M_{p\bar{p}}(M_{\Lambda\bar{\Lambda}})$ mass bins from the kinematic threshold to 4.5 GeV/c². The results of η_c are the mass of $M_{\eta_c} = 2.971 \pm 0.003^{+0.002}_{-0.001}$ GeV/c²(2.974±0.007^{+0.002}_{-0.001} GeV/c²) and the width of $\Gamma(\eta_c) = 48^{+8}_{-7} \pm 5$

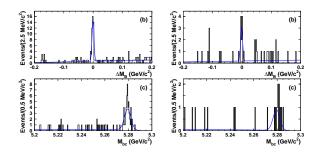


Figure 4: Candidate (a,b) $B^+ \to \Lambda_c^+ \Lambda_c^- K^+$ and (c,d) $B^0 \to \Lambda_c^+ \Lambda_c^- K^0$ decay events: (a,c) ΔM_B distribution for $M_{\rm bc} > 5.27~{\rm GeV/c^2}$ and (b,d) $M_{\rm bc}$ distribution for $|\Delta M_B| < 0.015~{\rm GeV/c^2}$. Curves indicate the fit results.

Modes	Yield	Efficiency(%)	$\mathcal{B}(J/\psi, \eta_c \to p\bar{p}, \Lambda\bar{\Lambda}) \times 10^{-3}$
$B^+ \to \eta_c K^+, \eta_c \to p\bar{p}$	$195.1^{+15.7}_{-14.7}$	35.8 ± 0.3	$1.58 \pm 0.12 ^{~+0.18}_{~-0.22} \pm 0.47$
$B^+ \to \eta_c K^+, \eta_c \to \Lambda \bar{\Lambda}$	$19.5^{+5.2}_{-4.5}$	5.3 ± 0.1	$0.87^{+0.24+0.09}_{-0.21-0.14} \pm 0.27$
$B^+ \to J/\psi K^+, J/\psi \to p\bar{p}$	$317.2^{+19.0}_{-18.0}$	37.3 ± 0.4	$2.21 \pm 0.13 \pm 0.31 \pm 0.10$
$B^+ \rightarrow I/2/2K^+ I/2/2 \rightarrow \Lambda \bar{\Lambda}$	$45.9^{+7.7}$	5.9 ± 0.3	$2.00^{+0.34}_{-0.32} + 0.34 + 0.08$

Table 1: Measured Branching Fractions $\mathcal{B}(J/\psi, \eta_c \to p\bar{p}, \Lambda\bar{\Lambda})$.

MeV/c²(40± 19 ± 5 MeV/c²) from $\eta_c \to p\bar{p}(\eta_c \to \Lambda\bar{\Lambda})$ mode. We define the η_c signal region as 2.940 GeV/c² < $M_{\Lambda\bar{\Lambda}}$ < 3.020 GeV/c². The fitted B signal yield, efficiency and branching fraction are shown in Table 1. In this study the decay $\eta_c \to \Lambda\bar{\Lambda}$ has been observed for the first time, with $\mathcal{B}(\eta_c \to \Lambda\bar{\Lambda}) = (0.87^{+0.24+0.09}_{-0.21-0.14} \pm 0.27) \times 10^{-3}$. The observed $\mathcal{B}(\eta_c \to \Lambda\bar{\Lambda})/\mathcal{B}(\eta_c \to p\bar{p})$ is $0.67^{+0.19}_{-0.16} \pm 0.12$ which is consistent with theoretical expectation⁶. We define the J/ψ signal region as 3.075 GeV/c² < $M_{p\bar{p}}(M_{\Lambda\bar{\Lambda}}) < 3.117$ GeV/c² and use events in this signal region to study the proton angular distribution in the helicity frame of the J/ψ . The helicity angle θ_X is defined as the angle between the proton flight direction and the direction opposite to the flight of the kaon in the J/ψ rest frame. The angular distribution of the kaon direction in the J/ψ rest frame is parameterized as $P(\alpha_B, \cos\theta_X) = (1 + \alpha_B \cos^2\theta_X)/(2 + 2\alpha_B/3)$ with $\alpha_B = (-2\alpha)/(\alpha+1)$. We determine $\frac{11}{\alpha_B}$ to be -0.60 \pm 0.13 \pm 0.14 $(p\bar{p})$ and -0.44 \pm 0.51 \pm 0.31 $(\Lambda\bar{\Lambda})$.

References

- 1. Belle Collaboration, M.-Z. Wang *et al.*, Phys. Rev. Lett. **90**, 201802(2003).
- 2. Belle Collaboration, M.-Z. Wang et al., Phys. Rev. Lett. **92**, 131801(2004).
- 3. Belle Collaboration, Y.-J. Lee, M.-Z. Wang et al., Phys. Rev. Lett. 93, 211801 (2004).
- 4. Belle Collaboration, M.-Z. Wang et al., Phys. Lett. B 617, 141(2005).
- 5. Belle Collaboration, Y.-J. Lee, M.-Z. Wang et al., Phys. Rev. Lett. 95, 061802(2005).
- 6. M. Anselmino, F. Caruso, S. Forte and B. Pire, Phys. Rev. D 38, 3516 (1988).
- 7. S. Kurokawa and E. Kikutani, Nucl. Instrum. Methods Phys. Res. A499, 1(2003).
- 8. Belle Collaboration, R. Chistov et al., hep-ex/0510074.
- 9. Belle Collaboration, N. Gabyshev et al., hep-ex/0508015.
- Belle Collaboration, T. Lesiak et al., Phys. Lett. B 605, 237 (2005), Erratum 617, 198 (2005).
- 11. F. Murgia and M. Melis, Phys. Rev. D **51**, 3487 (1995).